

Oscillating dark matter: experiments and data analysis at SYRTE

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SYRTE seminar, July 1st, 2021

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with valuable help from FOP + FMO teams and electronics + IT services

Menu

- Dark matter (DM) overview
- Ultralight dark matter and time/frequency metrology
- **DM interacting with standard model spins**
- **Non-universally coupled scalar field, variation of constants**
- **Atomic spectroscopy and optical cavities**
- **Experiments at SYRTE**
- Towards axion searches
- Conclusion

[Hees, Guéna, Abgrall, Bize, Wolf, PRL **117**, 061301, 2016]

[Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD **98**, 064051, 2018]

[Alonso, Blas, Wolf, JHEP **69**, 2019]

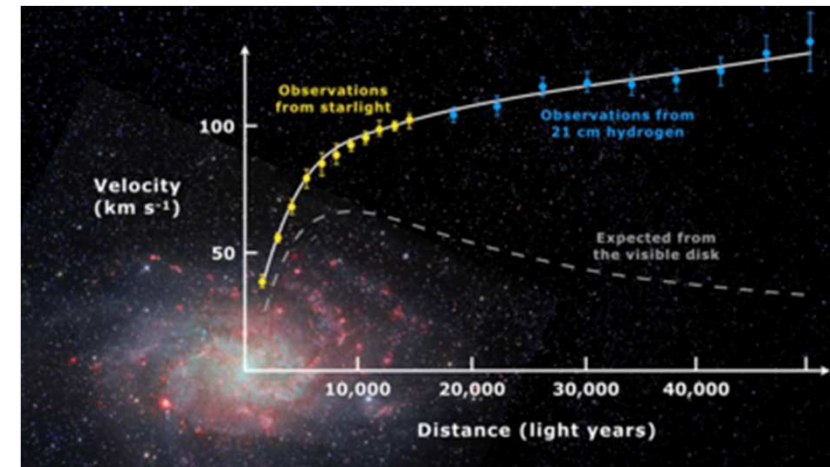
[Wolf, Alonso, Blas, PRD **99**, 095019, 2019]

[Roberts, et al., NJP **22**, 093010, 2020]

[Savalle E. et al., PRL **126**, 051301, 2021]

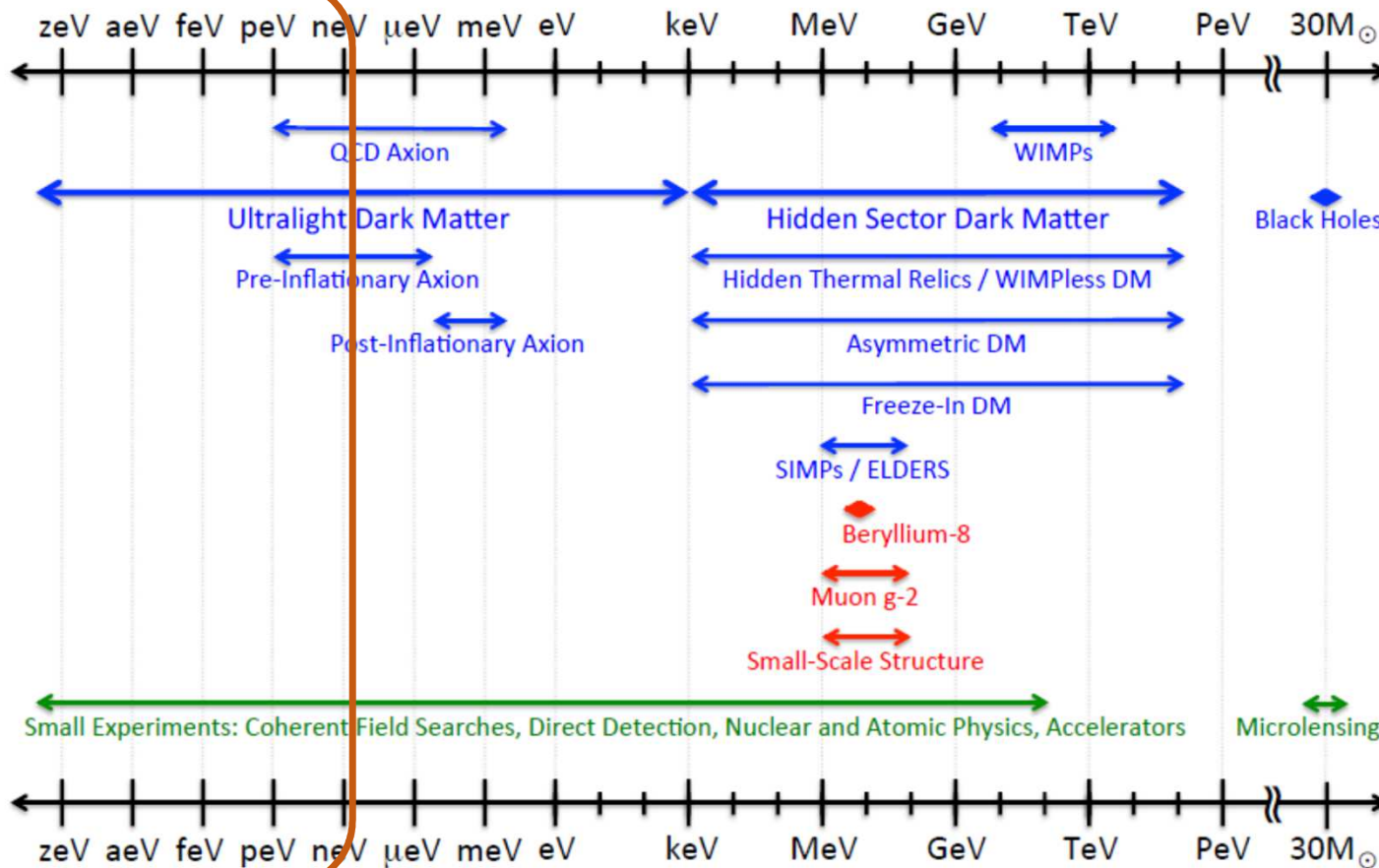
Dark Matter (1)

- “Evidence” for DM is purely gravitational, but of several types e.g.:
 - Galaxy rotation curves
 - Gravitational lensing
 - Cosmic Microwave Background
 - Structure formation
 - ...
- We “know” that:
 - It is cold ($v \ll c$)
 - Forms a galactic halo
 - Has virialized in the galaxy ($\delta v \approx 10^{-3} c$, $\langle v \rangle \approx 0$)
 - It’s energy density in the solar system is $\approx 0.4 \text{ GeV/cm}^3$ and $\langle v \rangle \approx 10^{-3} c$
- We hope that:
 - More gravitational evidence will be obtained to constrain its properties
 - DM interacts other than gravitationally with standard model fields
 - Someone will detect it locally
 - New physics will be learned



Dark Matter (2)

Dark Sector Candidates, Anomalies, and Search Techniques



From "US Cosmic Visions: New Ideas in Dark Matter 2017 : Community Report" arXiv:1707.04591

- Spans 90 orders of magnitude in mass !
- Here we will concentrate on low masses.
- In that region standard collisional (recoil based) detection techniques fail.

Ultralight DM

$$N_{occ} = \frac{n}{n_{\delta v}} \simeq \frac{3\pi^2 \hbar^3 \rho}{4m^4 \delta v^3}$$

$N_{occ} = 1$ in our galaxy for $m \approx 10$ eV.

- For $N_{occ} > 1$ the DM field can be treated as a classical field.
- It is likely to oscillate at its Compton frequency $\omega = mc^2/\hbar$.
- It may form “clumps” e.g. topological defects or relaxion stars/halos.
- For $N_{occ} < 1$ it must be quantized i.e. treated as a particle.
- Fermions cannot have $N_{occ} > g$ (g = number of internal degrees of freedom).
- Fermionic DM mass must be $> eV$.
- Bosonic DM can be treated as a classical field for mass below 10 eV or so.

Observable effects

1. DM fields interacting with the spin of the electrons or nuclei in the atoms.
 - ⇒ Effect on spin dependent atomic transition frequencies (Hyperfine transitions, Zeeman states, ...).

2. DM scalar field with non-universal scalar couplings to SM fields.
 - ⇒ Apparent violations of the equivalence principle
 - ⇒ Space-time variation of fundamental constants
 - ⇒ Change of atomic transition frequencies
 - ⇒ Change of Bohr-radius = length of solids

Non-universally coupled scalar fields

[Damour & Donoghue 2010]
[Stadnik & Flambaum 2014,2015]

$$S = \frac{1}{c} \int d^4x \frac{\sqrt{-g}}{2\kappa} [R - 2g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - V(\varphi)] \\ + \frac{1}{c} \int d^4x \sqrt{-g} [\mathcal{L}_{\text{SM}}(g_{\mu\nu}, \Psi) + \mathcal{L}_{\text{int}}(g_{\mu\nu}, \varphi, \Psi)]$$

$$\mathcal{L}_{\text{int}} = \frac{\varphi^i}{i} \left[\frac{d_e^{(i)}}{4\mu_0} F^2 - \frac{d_g^{(i)} \beta_g}{2g_3} (F^A)^2 \right. \\ \left. - c^2 \sum_{w=e,u,d} (d_{m_w}^{(i)} + \gamma_{m_w} d_g^{(i)}) m_w \bar{\psi}_i \psi_i \right]$$

- $i = 1, 2$ for linear or quadratic coupling
- With five dimensionless coupling constants $d_x^{(i)}$

Variation of constants

$$\mathcal{L}_{\text{eff}}^{\text{EM}} = \underbrace{-\frac{e^2 c}{16\pi\hbar\alpha} F^2}_{\text{Electromagnetism from Standard Model}} + \underbrace{+d_e\varphi \frac{e^2 c}{16\pi\hbar\alpha} F^2}_{\text{Electromagnetism from scalar field}} \simeq \frac{-e^2 c}{16\pi\hbar\alpha(1+d_e\varphi)} F^2$$

$$\alpha_{EM}(\varphi) = \alpha_{EM} \left(1 + d_e^{(i)} \frac{\varphi^i}{i}\right)$$

$$m_w(\varphi) = m_w \left(1 + d_{m_w}^{(i)} \frac{\varphi^i}{i}\right)$$

$$\Lambda_3(\varphi) = \Lambda_3 \left(1 + d_g^{(i)} \frac{\varphi^i}{i}\right)$$

$i = 1, 2$

- Fundamental constants (α , Λ_3 , m_i) are functions of φ , and vary if φ varies.
- Different atomic transitions depend differently on fundamental constants and thus their relative frequency varies with φ .
- The length of solids (e.g. optical cavities) is proportional to the Bohr radius ($\propto 1/(m_e\alpha)$) and thus varies with φ .
- Light speed is unchanged (in geometric optics approximation)

[Damour & Donoghue 2010]

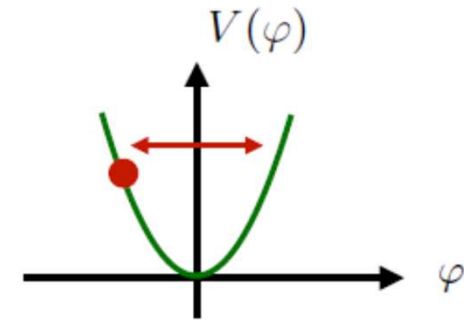
[Stadnik & Flambaum 2014,2015]

Evolution of the galactic scalar field (2)

$$V(\varphi) = 2\frac{c^2}{\hbar^2}m_\varphi^2\varphi^2$$

- Assume a quadratic potential for φ .
- Varying Lagrangian with respect to φ gives a KG equation:

$$\frac{1}{c^2}\ddot{\varphi}(t, \mathbf{x}) - \Delta\varphi(t, \mathbf{x}) = -\frac{4\pi G}{c^2}f(d_j^{(i)})\rho_A(\mathbf{x})\varphi(t, \mathbf{x})^{i-1} - \frac{c^2m_\varphi^2}{\hbar^2}\varphi(t, \mathbf{x})$$



$$\varphi^{(1)}(t, \mathbf{x}) = \varphi_0 \cos(\omega t + \delta) - s_A^{(1)} \frac{GM_A}{c^2 r} e^{-r/\lambda_\varphi}$$

$$\varphi^{(2)}(t, \mathbf{x}) = \varphi_0 \cos(\omega t + \delta) \left[1 - s_A^{(2)} \frac{GM_A}{c^2 r} \right]$$

- The solutions for a spherically symmetric mass distribution oscillate at $\omega = m_\varphi c^2 / \hbar$
- The Yukawa term has range $\lambda_\varphi = \hbar / (m_\varphi c)$
- s_A are functions of d_j and the central body (GM_A/R_A)
- Linear ($i=1$) solution is well known
- Quadratic ($i=2$) solution is less common and has interesting phenomenology

[Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD **98**, 064051, 2018]

Link to Dark Matter

$$\rho_{\bar{\varphi}} = \frac{c^2}{4\pi G} \frac{\omega^2 \varphi_0^2}{2} = \frac{c^6}{4\pi G \hbar^2} \frac{m_\varphi^2 \varphi_0^2}{2}$$

- The cosmological density (+) and pressure (-) of φ are given by $\frac{c^2}{8\pi G} \left(\dot{\varphi}^2 \pm \frac{V(\varphi)c^2}{2} \right)$.
- The oscillating part of $\varphi(t)$ has zero average pressure and is therefore a candidate for Dark Matter.
- Equating its average density at spatial infinity with the DM density ($\approx 0.4 \text{ GeV/cm}^3$) fixes the amplitude φ_0 .
- The oscillation translates into an oscillation of the fundamental constants that can be searched for in a 6 parameter space (m_φ, d_x).
- The mass m_φ is given by the frequency of oscillation, the coupling constants d_x by the amplitude.

[Stadnik & Flambaum 2014, 2015]

[Arvinalaki, Huang, Van Tilburg 2015]

Evolution of the galactic scalar field (2)

Coherence time:

$$\hbar\omega = mc^2 + \frac{mv^2}{2} \Rightarrow \frac{\delta\omega}{\omega} \approx \frac{v\delta v}{c^2} \approx 10^{-6}$$

for $\delta v \approx v \approx 10^{-3} c$ in the virialized galaxy

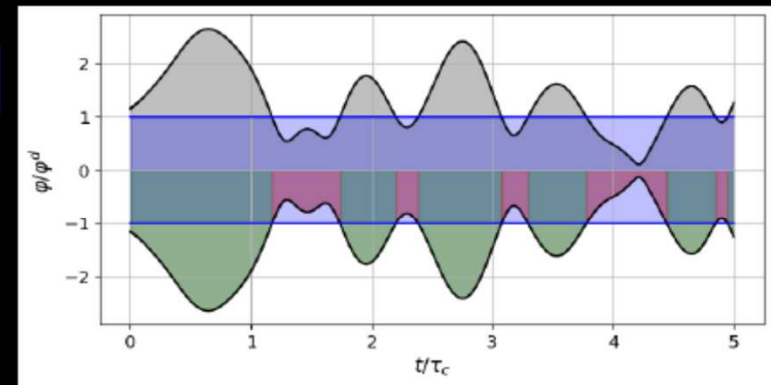
$$\delta\omega \tau_{coh} = 2\pi$$

- In our Cs/Rb experiment [Hees et al. 2016] ($f < 5.7 \times 10^{-4}$ Hz) this gives $\tau_{coh} > 55$ years.
 - In the DAMNED experiment [Savalle et al. 2021] ($f = [10:200]$ kHz) this gives $\tau_{coh} = [5:100]$ s.
- The velocity distribution is stochastic and that needs to be taken into account either by decreased sensitivity [Centers et al. arXiv:1905.13650] or by modelling the full stochastic evolution.

Stochastic scalar field

$$\varphi = \varphi_0 \sum_{j=1}^{N_j} \alpha_j \sqrt{f_{DM}(\omega_{\varphi_j}) \Delta\omega} \cos(\omega_{\varphi_j} t + \delta_j)$$

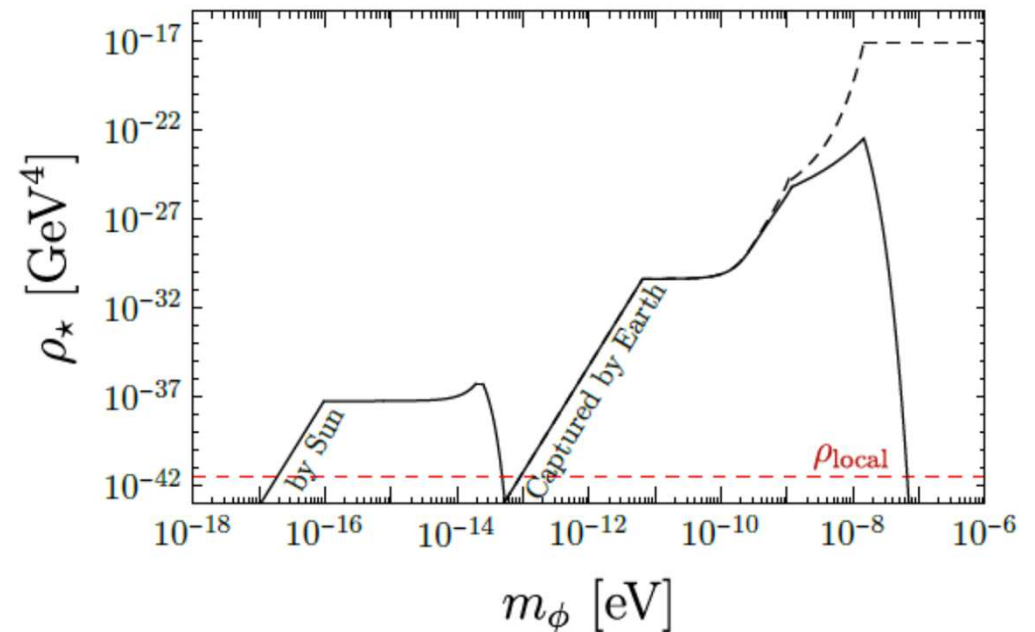
J. W. Foster et al. - PRD (2018)



SYRTE

Scalar field transient “clumps” (2)

- Other self potentials than quadratic are possible.
- The scalar field may form objects (boson stars) or halos around standard matter objects (e.g. Earth, Sun), or topological defects (e.g. domain walls)
- The resulting field may still oscillate at its Compton frequency ($\omega = m_\phi c^2 / \hbar$).
- This could lead to an overdensity around massive objects like the Earth, or to transient local variations of the scalar field.
- It may also modify the coherence properties of the field (e.g. much longer coherence time)



[Derevianko, A. & Pospelov, M., Nature Physics, **10**, 933, 2014]

[Banerjee, A.; Budker, D.; Eby, J.; Kim, H. & Perez, G., Communications Physics **3**, 1, 2020]

Dark activities at SYRTE

Theory:

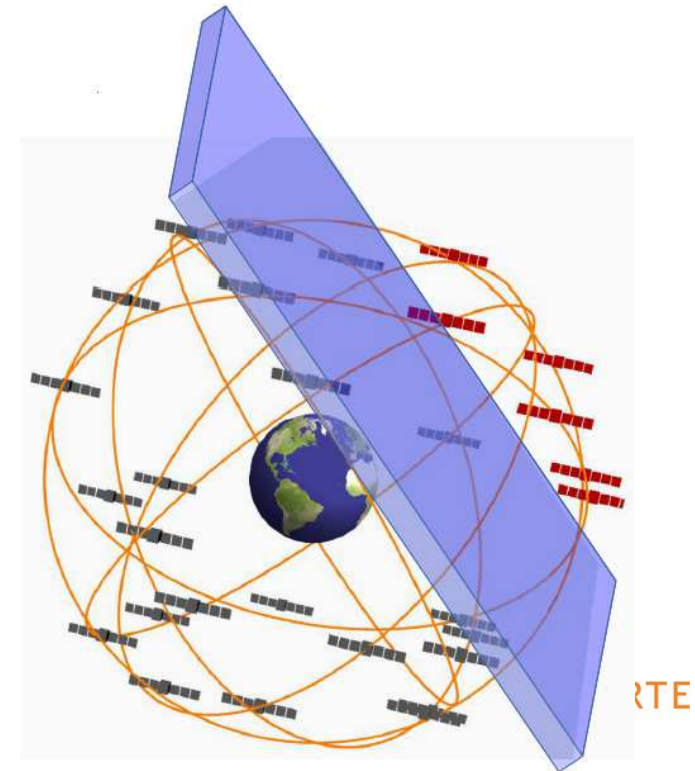
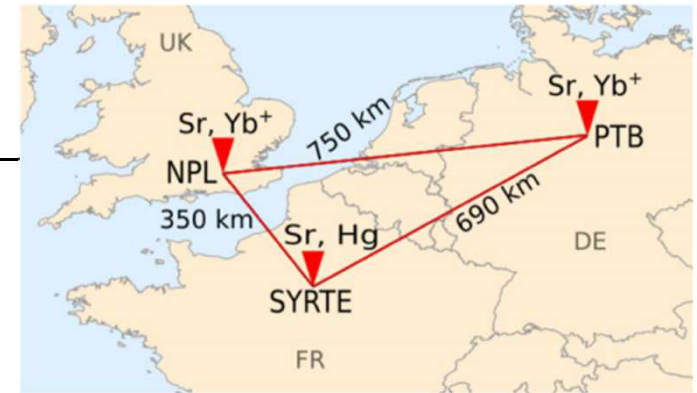
- Extensive study/review of equivalence principle violating scalar DM, scalar coupling, atomic transitions/free fall tests [Hees 2018]
- Study of interactions with atomic spins: scalar/fermion/vector boson DM, with axial/tensor coupling, contact interaction/mediator. Effect in atomic clocks and co-magnetometers [Alonso 2019, Wolf 2019].

Experiments:

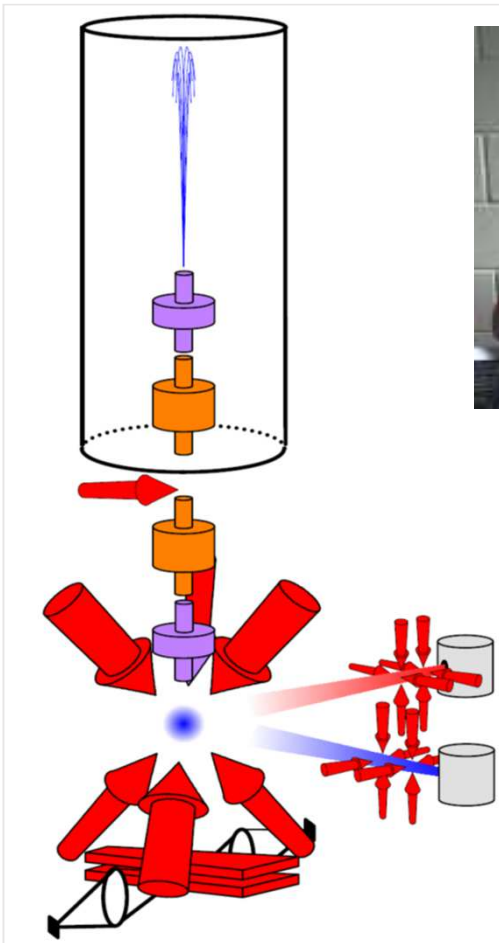
- Rb/Cs dual cold atom clock, long term comparison [Hees 2016]
- The DAMNED experiment [Savalle 2021]
- Europe-wide comparison of optical clocks to search for transients [Roberts 2020]
- The GASTON project (GALileo Survey of Transient Objects Network), searching for transients using the clocks on board the Galileo satellite constellation [ESA contract, ongoing].

Collaborations:

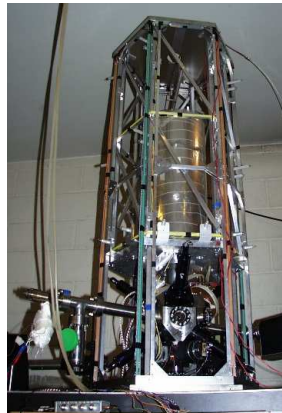
- Collaboration with CEA-IRFU, SHUKET and G-LEAD experiments.
- Common monthly seminars with Univ. Tokyo DM team.



The SYRTE dual Rb-Cs fountain FO2



André Clairon
1947 - 2015



- Built in early 2000s by André Clairon and co-workers.
- Operates simultaneously on laser cooled (μK) ^{87}Rb and ^{133}Cs since 2008 (common mode systematics).
- Most accurate and stable Rb/Cs frequency ratio measurement world-wide (and longest duration).
- Contributes continuously to TAI with both Rb and Cs
- Previously used to constrain linear drifts of fundamental constants, and variations proportional to U/c^2 i.e. annual variations [Guéna, PRL 2012]+updates.
- All systematics are evaluated and corrected during operation.

[Guéna et al. 2010, 2012, 2014]

Atomic Spectroscopy

- Different atomic transition frequencies depend differently on fundamental constants.
- Comparison of two atomic transition frequencies ($Y=X_A/X_B$) is a direct measure of the scalar field. Can be used to search for the space-time variation of $\varphi(t, \mathbf{x})$.

$$\frac{Y(t, \mathbf{x})}{Y_0} = K + \left(\kappa_{X_A}^{(i)} - \kappa_{X_B}^{(i)} \right) \varphi^i(t, \mathbf{x})$$

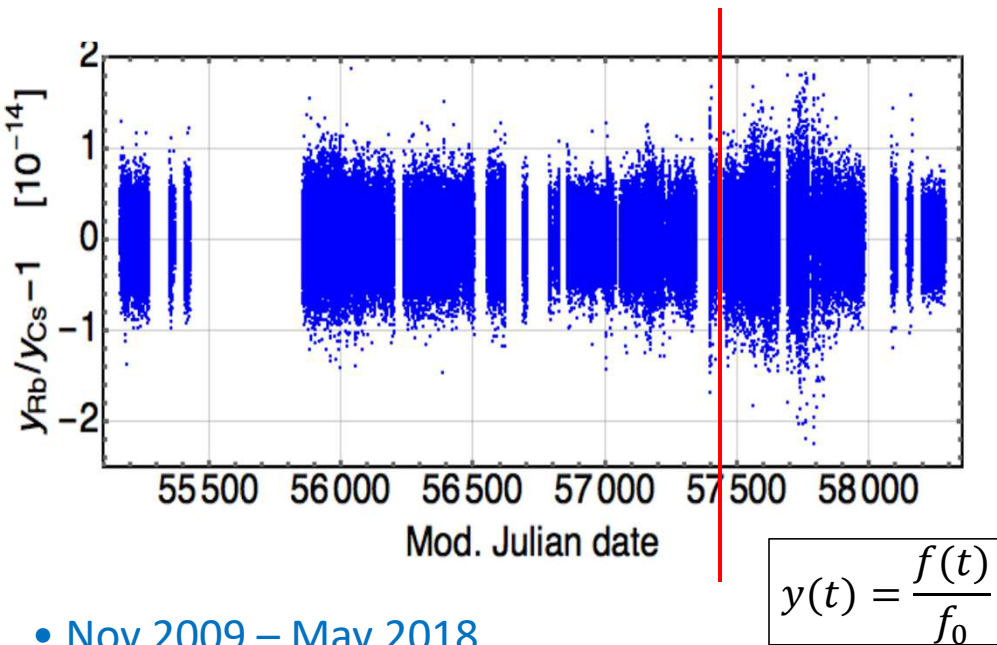
The sensitivity coefficients $\kappa_X^{(i)}$ involve the coupling constants $d_j^{(i)}$ and are obtained from atomic and nuclear structure calculations (Flambaum and co-workers [2006, 2008, 2009]).

$$\varphi^{(1)}(t, \mathbf{x}) = \varphi_0 \cos(\omega t + \delta) - s_A^{(1)} \frac{GM_A}{c^2 r} e^{-r/\lambda_\varphi}$$

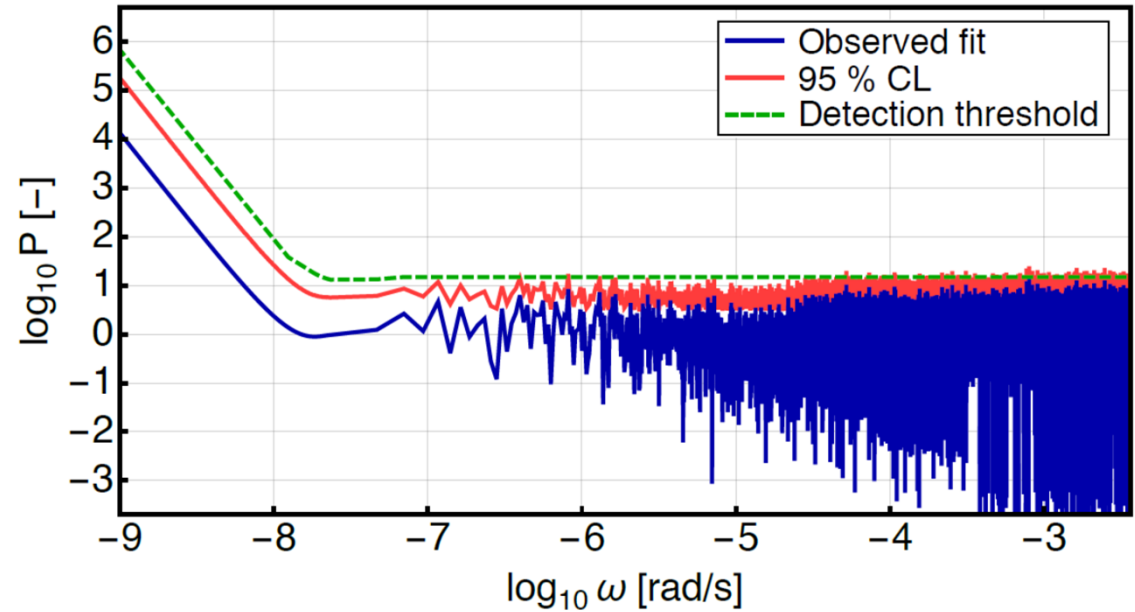
$$\varphi^{(2)}(t, \mathbf{x}) = \varphi_0 \cos(\omega t + \delta) \left[1 - s_A^{(2)} \frac{GM_A}{c^2 r} \right]$$

Can search for both, oscillations and spatial dependence in the field of body A (e.g. Earth)

FO2 Rb/Cs raw data



- Nov 2009 – May 2018
- Averaged to 100 points/day
- 144000 points in total
- $\approx 45\%$ duty cycle with gaps due to maintenance and investigation of systematics
- Standard deviation $\approx 3 \times 10^{-15}$

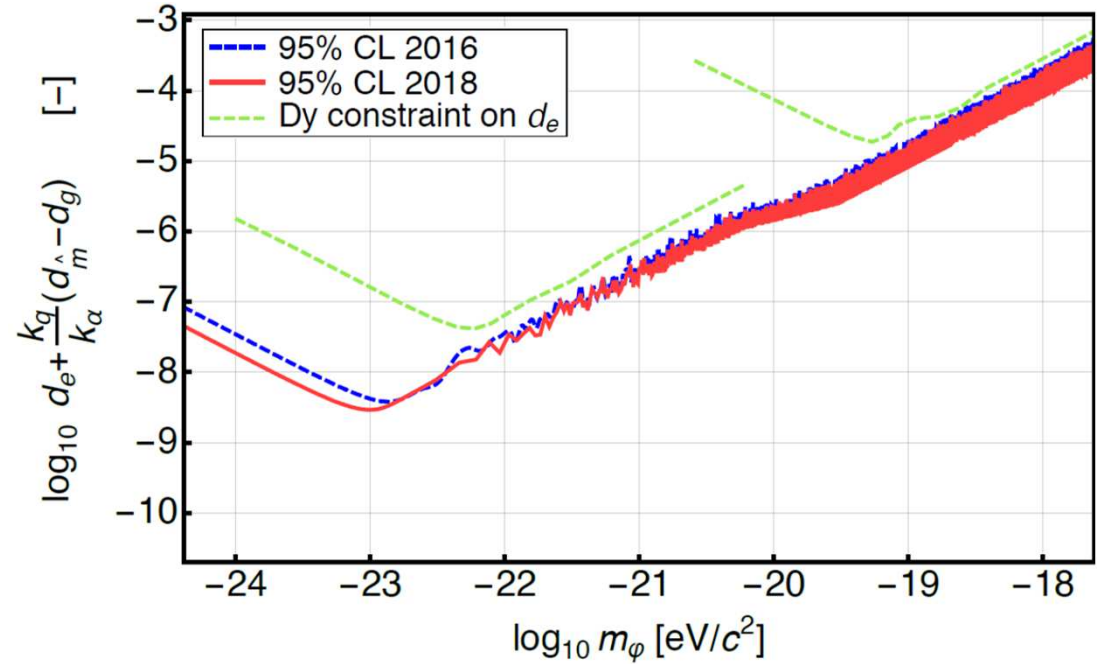
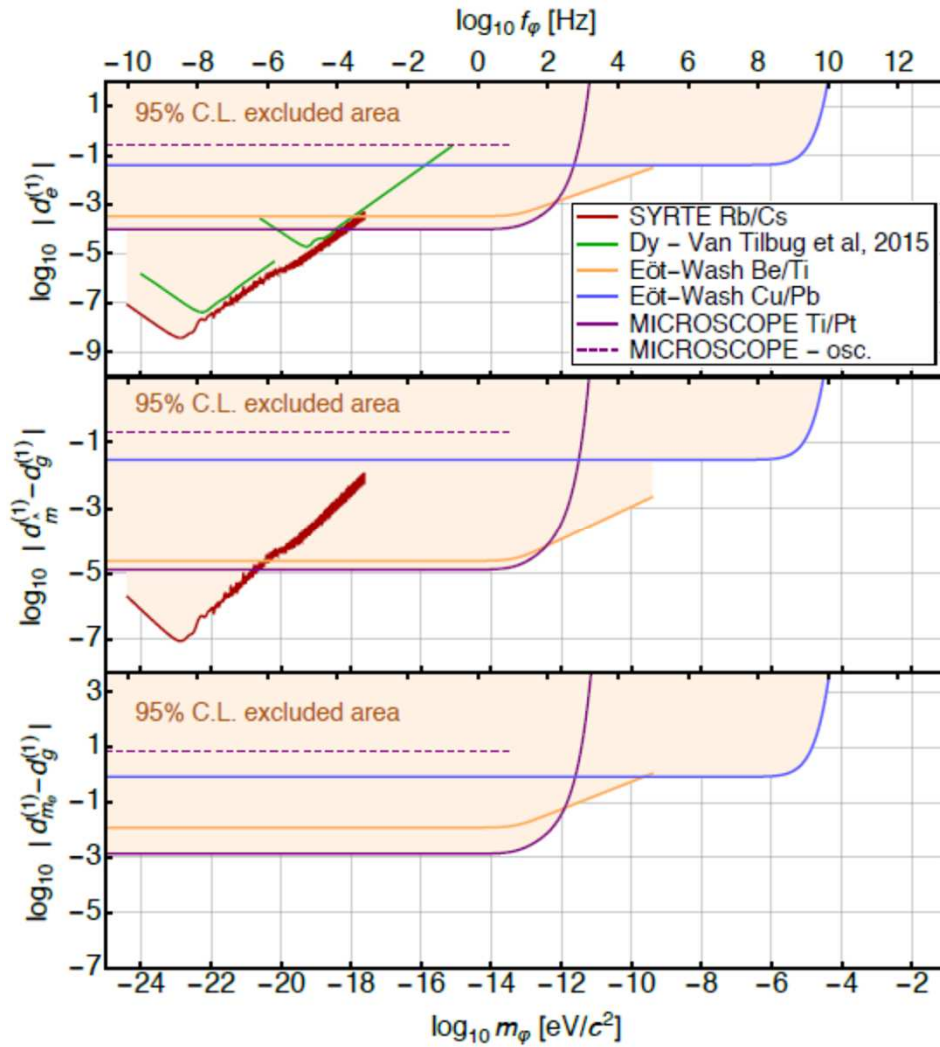


Search for a peak in normalized power

$$P_{\omega} = \frac{N}{4\sigma_0^2(\omega)} (C_{\omega}^2 + S_{\omega}^2)$$

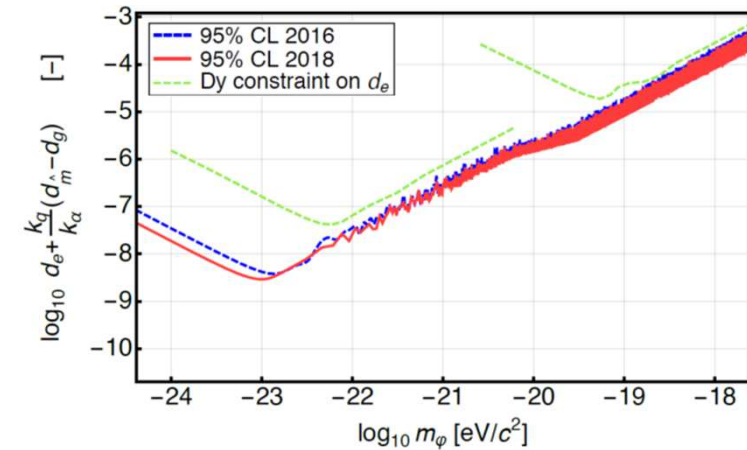
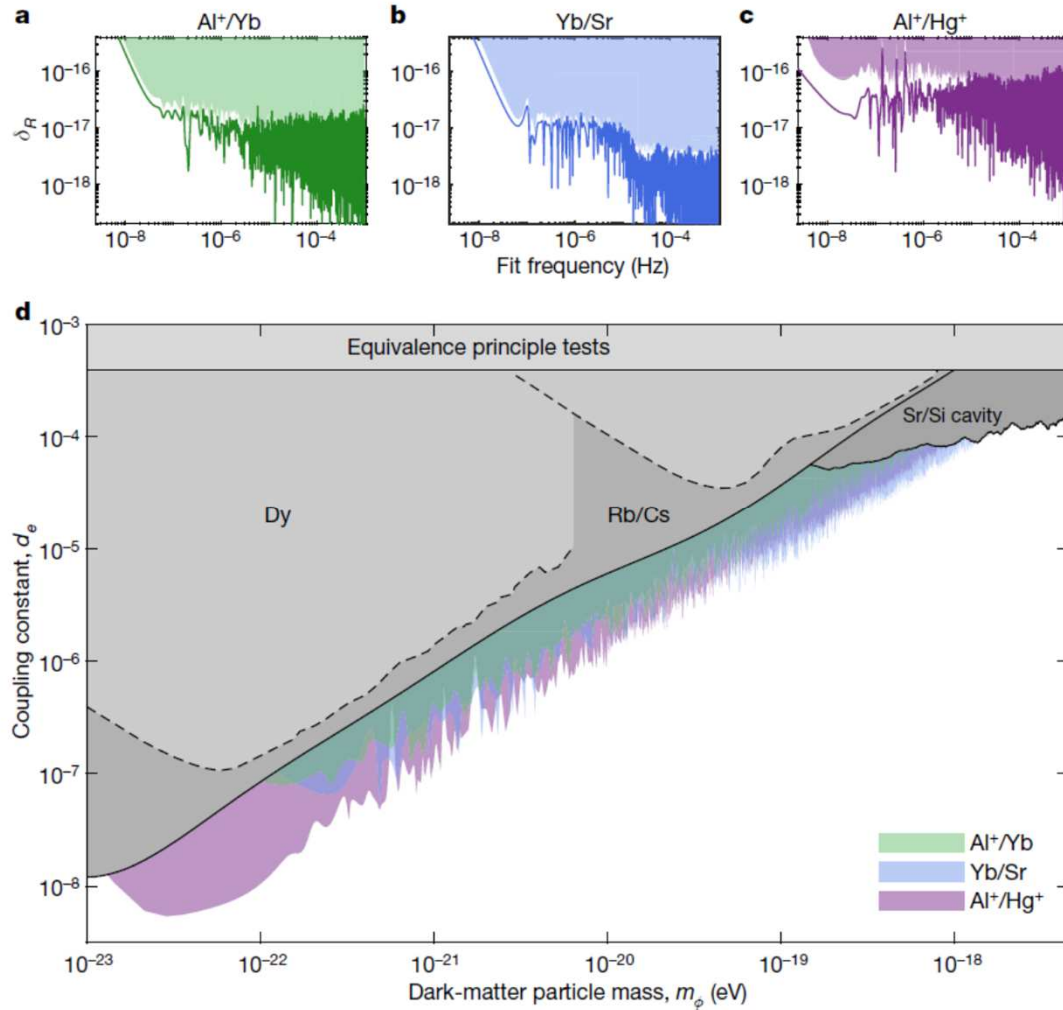
Update of [Hees, Guéna, Abgrall, Bize, Wolf, PRL **117**, 061301, 2016]

Constraints for linear coupling



[Hees 2016, Hees 2018]

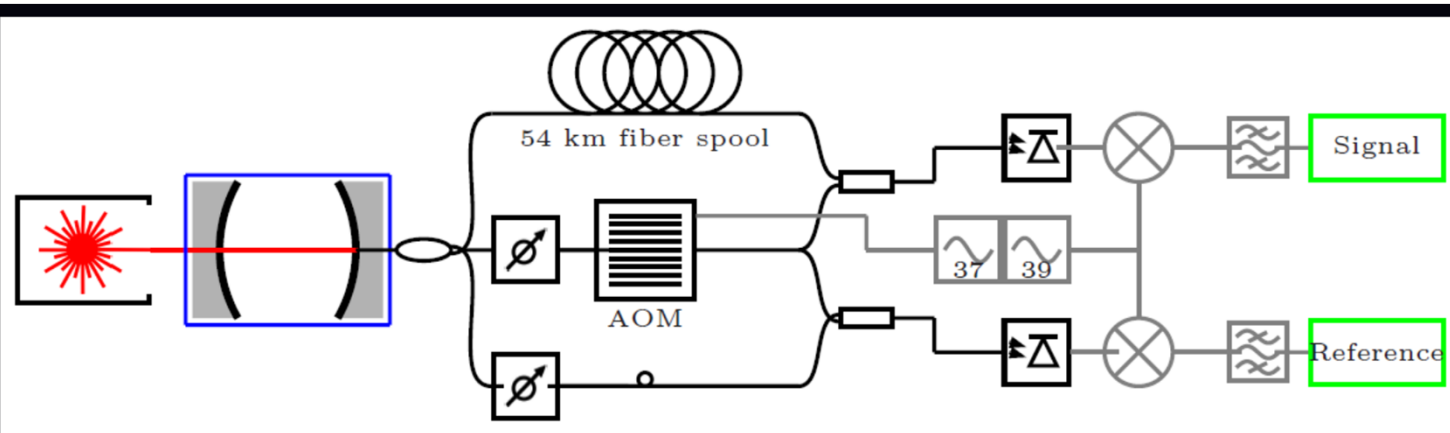
Constraints for linear coupling



- NIST-JILA optical clocks.
- All clock constraints « rescaled » to account for stochastic DM fluctuations.
- Improvement mostly due to α sensitivity.
- Optical clocks couple only to α but not to quark masses => continue using hyperfine transitions.

[Beloy et al., 2021]

DARk Matter from Non Equal Delays (DAMNED)



Source

1542 nm laser source stabilized on an ultra stable cavity, unevenly distributed in three arms.

Delay

Long delay line of 50 → 125 km
Short delay line of 1 m
AOM for detection

Detection

Beatnotes
Self-heterodyne
Photodiodes
Ettus X310

- Based on ultra-stable optical cavity and fibre delay line
- Variation of fundamental constants leads to variation of cavity/fibre length (Bohr radius) and fibre refractive index.
- Aiming at high frequency (10-100 kHz i.e. DM mass around 10^{-10} eV).

DAMNED signal

Cavity frequency oscillation

$$\omega(t) = \omega_0 + \Delta\omega(t) + \delta\omega \sin(\omega_\varphi t)$$

Color code

Nominal value

Noise

Dark matter effect

Fiber delay oscillation

$$T(t) = T_0 + \int_{t-T_0}^t \frac{\Delta T(t')}{T_0} dt' + \delta T \sin\left(\omega_\varphi t - \omega_\varphi \frac{T_0}{2}\right) \text{sinc}\left(\omega_\varphi \frac{T_0}{2}\right)$$

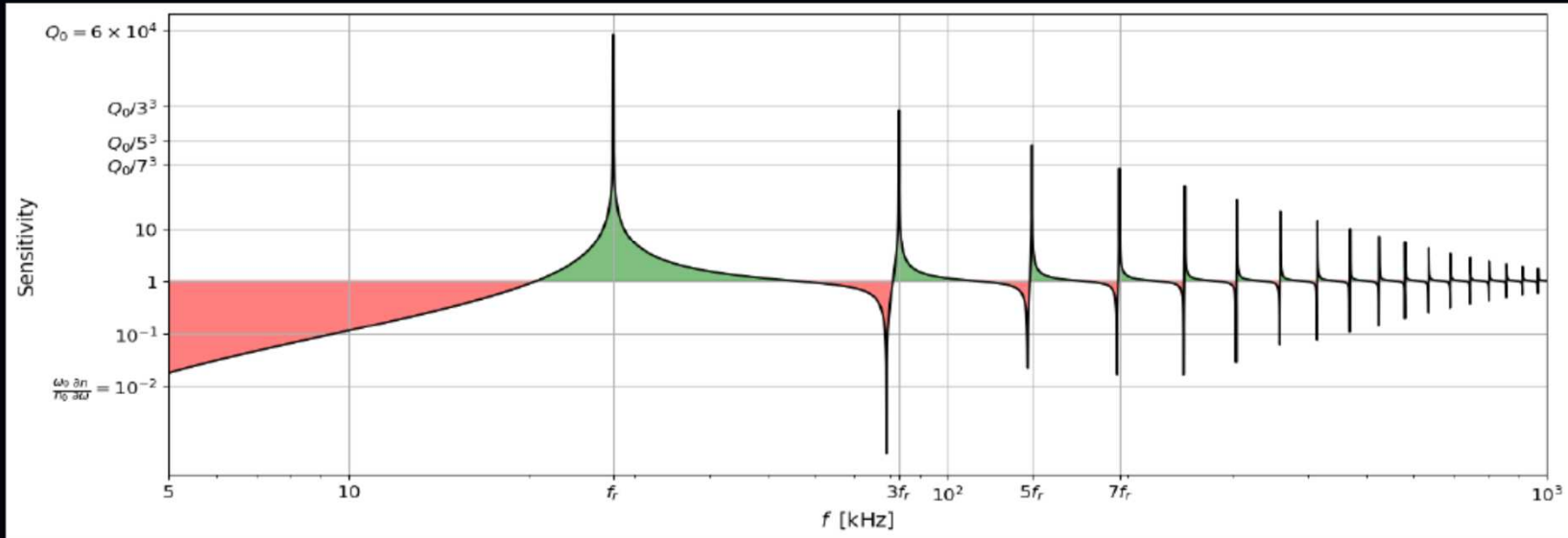
Phase difference between the delayed and non delayed signals

$$\begin{aligned} \Delta\Phi(t) = & \omega_0 T_0 + \omega_0 \int_{t-T_0}^t \left(\frac{\Delta T(t')}{T_0} + \frac{\Delta\omega(t')}{\omega_0} \right) dt' \\ & + \omega_0 T_0 \left(\frac{\delta T}{T_0} + \frac{\delta\omega}{\omega_0} \right) \sin\left(\omega_\varphi t - \omega_\varphi \frac{T_0}{2}\right) \text{sinc}\left(\omega_\varphi \frac{T_0}{2}\right) \end{aligned}$$

Use two fibre lengths (52 and 56 km) to avoid sensitivity zeros from sinc function.

DAMNED sensitivity

Full sensitivity



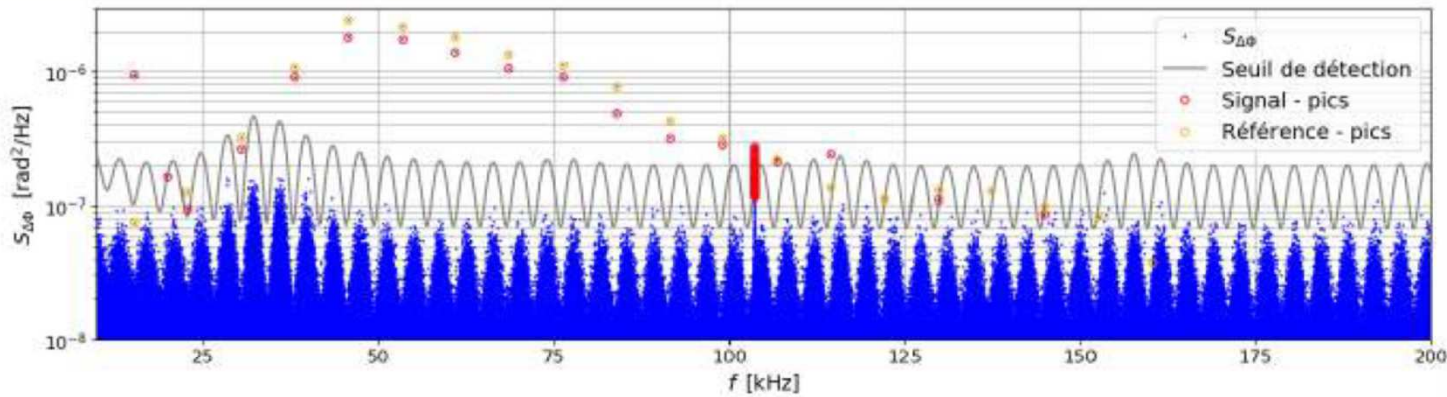
Best sensitivity at mechanical resonances of the cavity

Link to the coupling constants

$$\left(\frac{\delta\omega}{\omega_0} + \frac{\delta T}{T_0} \right) \simeq d_e \varphi_0 \text{ "Sensitivity"}$$
$$\text{or } \simeq d_{m_e} \varphi_0 \text{ "Sensitivity"}$$

DAMNED data

Mesure de la phase



Limites

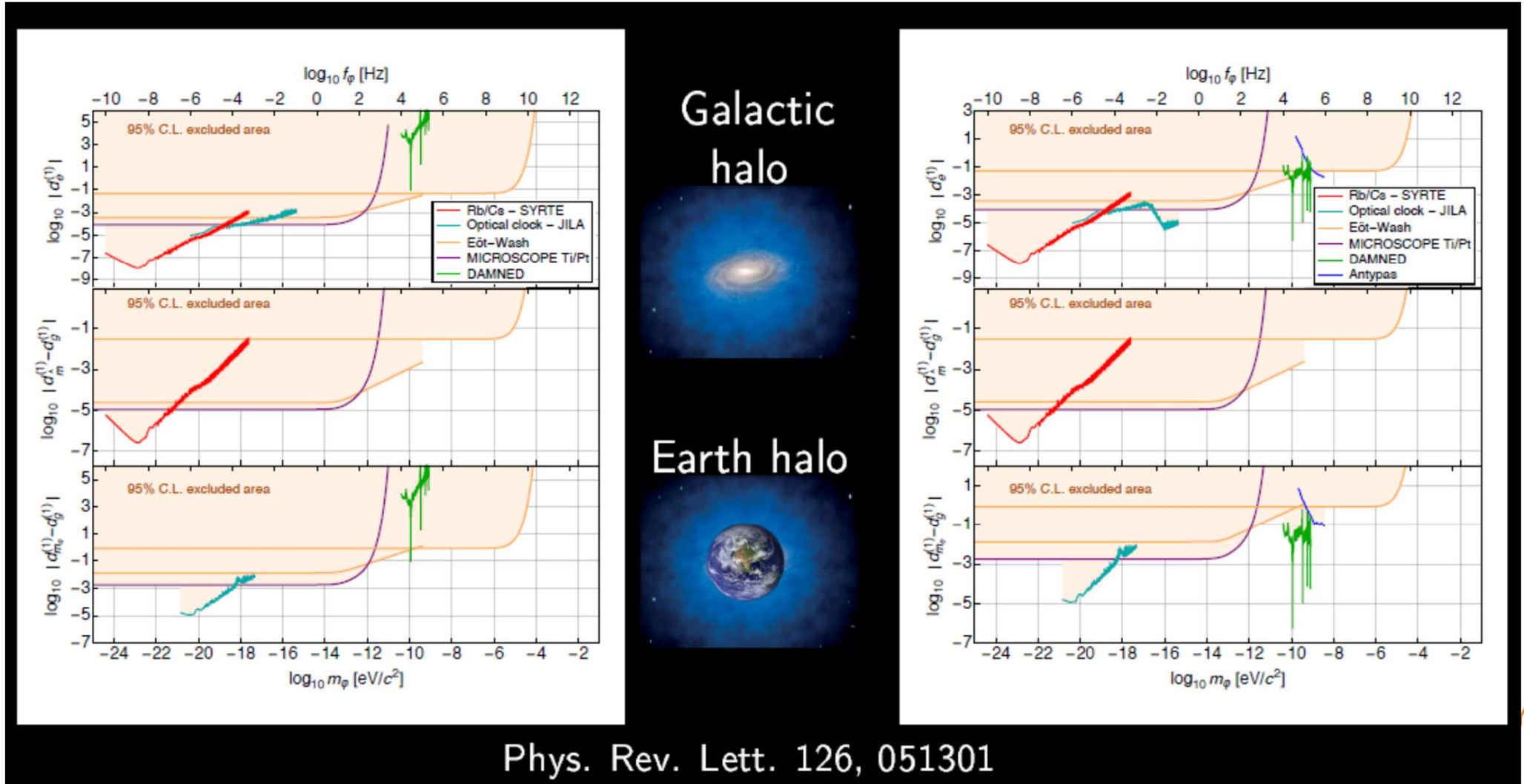
- < 10 kHz, bruit thermique/acoustique des fibres
- > 200 kHz, boucle de stabilisation laser/cavité

Fonction de transfert

$$S_{\Delta\phi}(\omega) = 4 \sin^2\left(\frac{\omega T}{2}\right) S_{\phi_c}$$

- 12 days data at 500 kHz sampling
- Limited by cavity noise after air conditioning failure
- Peaks excluded either by ref. arms or shape and e.g. temperature dependence or by using another cavity.
- Use stochastic Bayesian analysis (Derveianko, Hees)

DAMNED results



Towards Axion searches

- Axions and axion-like particles are “popular” pseudoscalar ultra-light DM candidates
 - Initially postulated to solve the “strong-CP” problem in QCD, aka the problem of vanishing electric dipole moment of the neutron
 - They have similar distributions as the scalar fields I discussed so far
- ⇒ Similar data analysis methods
- Generally searched for using strong magnetic fields to induce RF photon emissions. E.g. G-LEAD experiment @ CEA-IRFU.
 - They may also induce polarization dependent effects in cavities e.g. DANCE experiment @ Tokyo Univ., or [Goryachev 2019].

PhD thesis of Jordan Gué starting, hopefully, Oct. 2021 to theoretically study and model possible experimental axion searches using the optical cavities at SYRTE.

Conclusion

- Dark matter searches have become a major “application” of the outstanding uncertainty achieved in time/frequency metrology
- SYRTE has contributed significantly in that domain, and is continuing to do so (e.g. GASTON project).
- Unique data acquisition and analysis methods were developed for DAMNED, and are being adapted to other experiments (e.g. axion searches).
- There is still plenty of unexplored parameter space, may be accessible with SYRTE technology and expertise. And who knows what we will find?
- All of this activity relies on the interplay between theory and experimental teams in several groups and services of SYRTE. **Thanks to all who make an effort for fruitful collaborations!**